

D. DECAY PROCESSES

D.01. Decay modes

D.02. Decay law

D.03. Alpha decay & cluster emission

D.04. Beta decay

D.05. Gamma decay

D.06. Fission & fusion

D.01. Decay modes

α decay: emission of ${}^4\text{He}$

β decay: emission of e^- or e^+

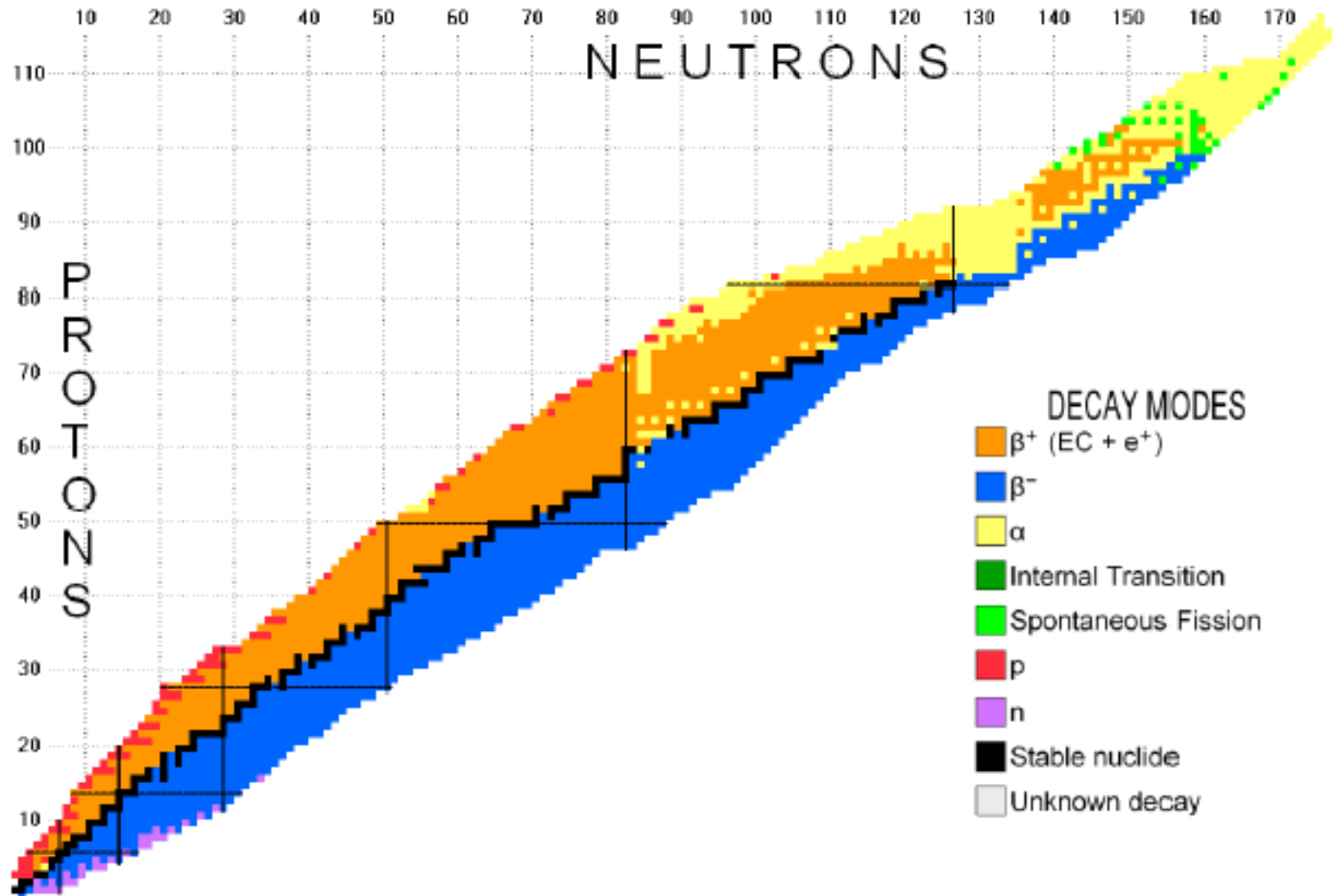
γ decay: emission of high energy photons

p decay: particle (proton/neutron) emission

c decay: emission of heavy clusters (C,O,Ne,Mg)

sf: spontaneous fission

Nuclear chart of decay modes



D.02. Decay law

$$dN \propto N$$

$$\frac{dN}{dt} = -\lambda N$$

$$N(t) = N_0 e^{-\lambda t}$$

λ = decay constant = probability of a nucleus decaying per second

Half-life = time for half the nuclei to decay

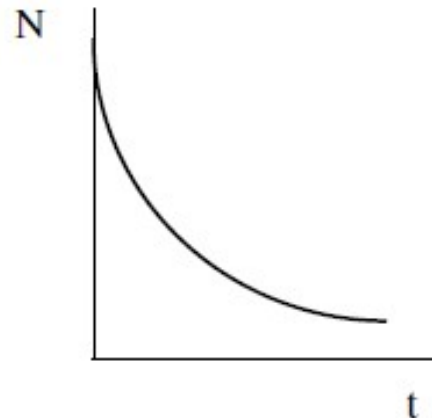
$$t_{1/2} = \frac{\ln 2}{\lambda}$$

Lifetime (average)

$$\tau = \frac{1}{\lambda}$$

Decay width is defined by:

$$\Gamma = \hbar \lambda$$



Wave function describing a decaying state

Decay law $|\Psi_{ext}(R, t)|^2 \rightarrow N(t) = N_0 e^{-\lambda t}$

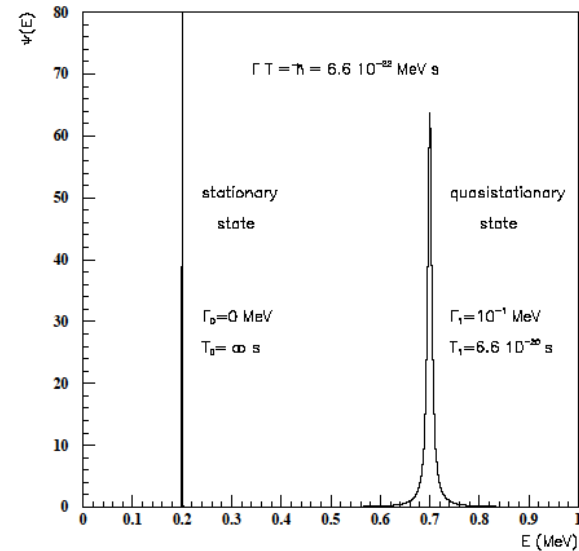
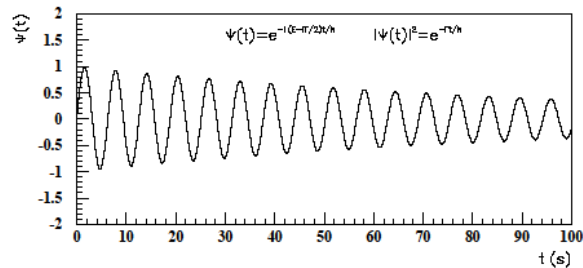
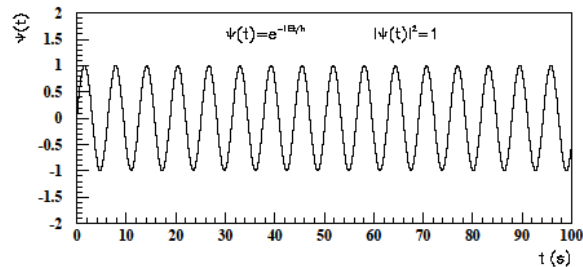
implies an external
wave function

$$\Psi_{ext}(R, t) = \Psi_{ext}(R) e^{-i(E/\hbar)t}$$

with complex
energy

$$E = E_0 - i \frac{\Gamma}{2}$$

The width of the energy distribution
 given by the Fourier transform of a decaying
 time-dependent signal
 is proportional to the decay width



**Decay rate (activity)
for the emission process**

P(arent) → D(aughter)+C(luster)

$$R = \left| \frac{dN}{dt} \right| = \lambda N_0 e^{-\lambda t} = \lambda N(t)$$

Units

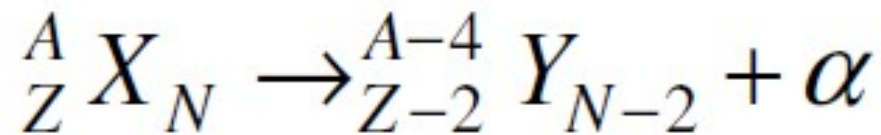
Becquerel (B) = 1 decay / second

Curie (C) = $3.7 \cdot 10^{10}$ B (1 gram of Radium)

D.03. Alpha decay

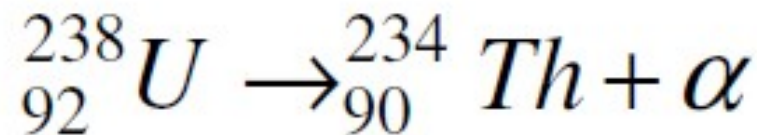
α -particle is a ${}^4\text{He}$ nucleus (2p & 2n)

- Z ↓ by 2
- N ↓ by 2
- A ↓ by 4



“parent”

“daughter”



$$t_{1/2} = 4.47 \times 10^9 \text{ yrs}$$

Disintegration energy

$$Q = (M_P - M_D - m_\alpha) c^2$$

$$M_U = 238.050784 \text{ u}$$

$$M_{Th} = 234.043593 \text{ u}$$

$$m_\alpha = 4.002602 \text{ u}$$

$$M_U - M_{Th} - m_\alpha = 0.004589 \text{ u}$$

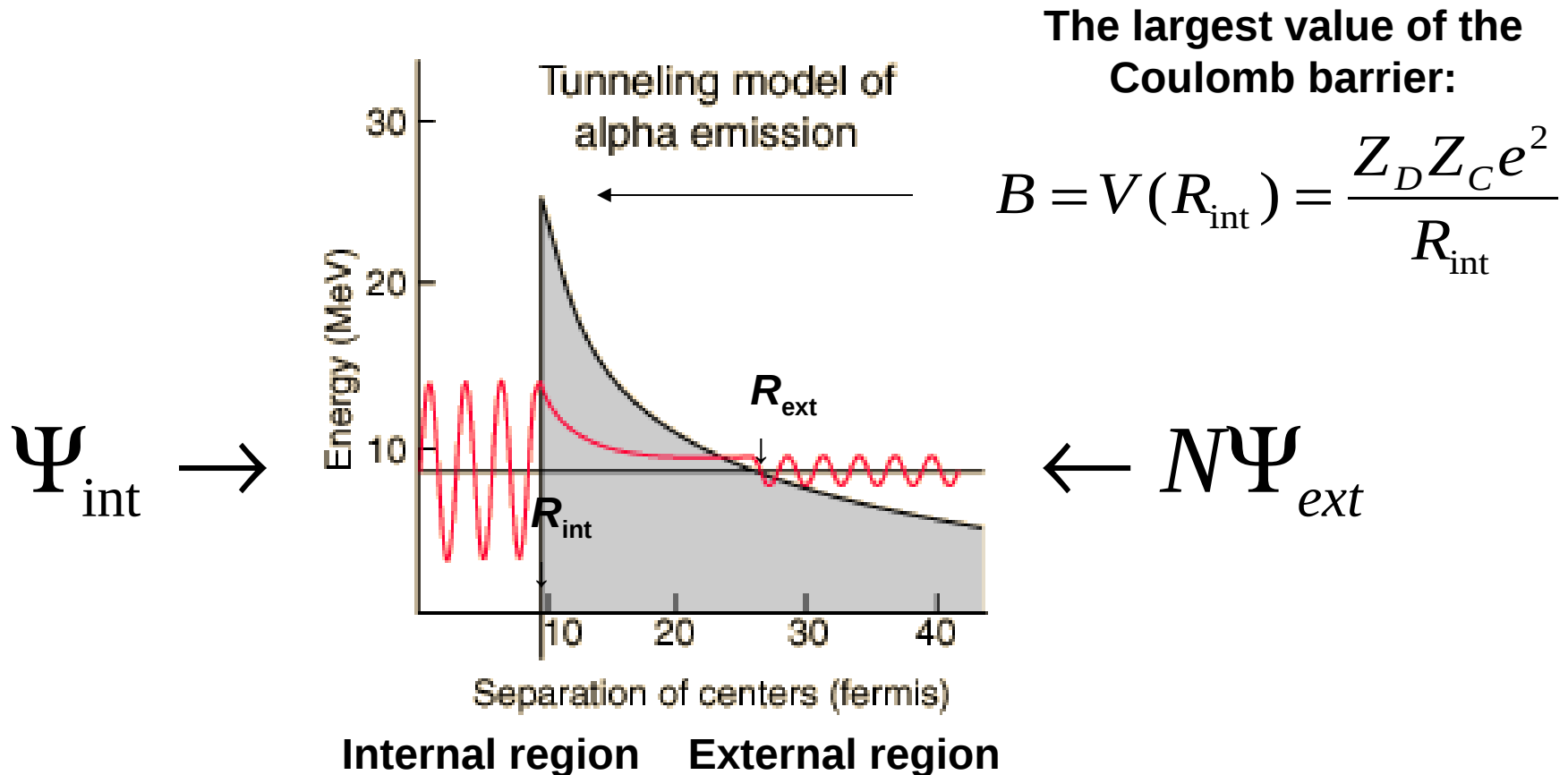
$$Q = \Delta m \cdot c^2 = 0.004589 \times 931.502$$

$$= 4.275 \text{ MeV}$$

Theory of the α -decay (and emission of heavier charged clusters)

The first probabilistic interpretation of the wave function for a particle penetrating a repulsive Coulomb barrier was given by

G. Gamow "Zur Quantentheorie des Atomkernes" (On the quantum theory of the atomic nucleus), *Zeitschrift für Physik*, vol. 51, 204-212 (1928).



**Geiger-Nuttall law
in terms of the Coulomb parameter**

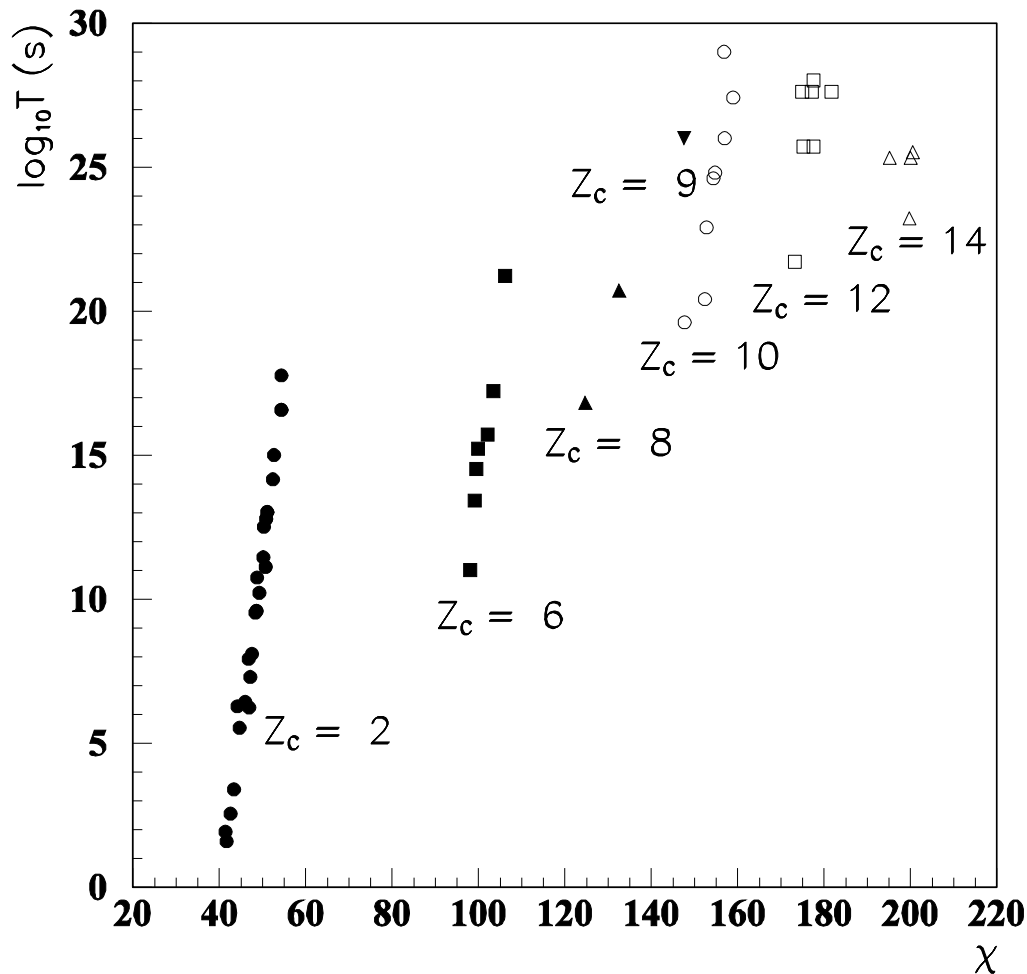
$$\log_{10} \Gamma = a\chi + b$$

where the Coulomb parameter is:

$$\chi = \frac{2Z_D Z_C}{\hbar v} \propto \frac{2Z_D Z_C}{\sqrt{E}}$$

$$Z_C = 2$$

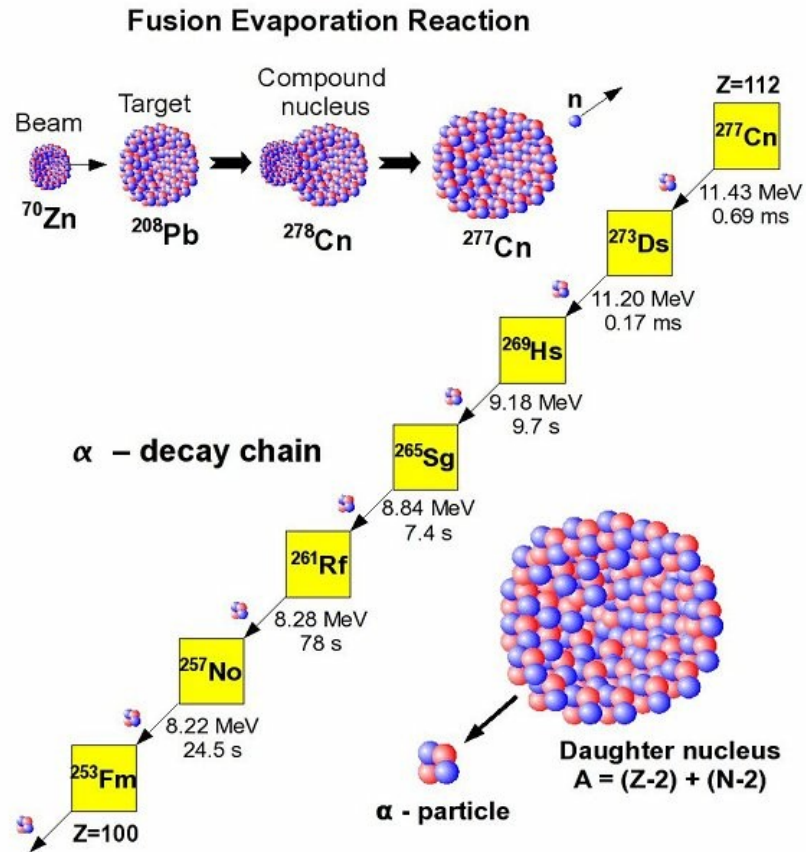
Quantum penetration explains Geiger-Nuttall law for α & cluster decays (emission of C, O, Ne, Mg, Si)



Coulomb parameter
In cluster emission:

$$\chi = \frac{2Z_D Z_C e^2}{\hbar v} \rightarrow c \frac{Z_D Z_C}{\sqrt{E}}$$

Superheavy nuclei are formed by fusion and detected by alpha decay chains



D.04. Beta decay

$$\beta^- : n \rightarrow p + e^- + \bar{\nu}_e$$

Decay of a free neutron. 'Lifetime' ~ 886 s

Electron-anti-neutrino

$$\beta^+ : p \rightarrow n + e^+ + \nu_e$$

Electron-neutrino

Cannot occur for a FREE proton ? 29

Neutrino

$$p \rightarrow n + e^+$$

$$\frac{1}{2} \quad \frac{1}{2} \quad \frac{1}{2}$$

- Three fermions: all have spin = $\frac{1}{2}$
- Spin angular momentum is not conserved in the above reaction
- The 'other' particle must have spin = $\frac{1}{2}$

Theory of β -decay

Describe the strength of the interaction
with an operator ('potential') V
Define a "matrix element"

$$V_{fi} = \int \psi_f^* V \psi_i dv$$

The 'initial' and 'final' states are
coupled by the operator V

Fermi's Golden Rule

The transition probability is given by

$$\lambda_{fi} = \frac{2\pi}{\hbar} |V_{fi}|^2 \rho_f$$

The nature of the interaction was
unknown at the time !

The density of states can be calculated
in 'Phase Space' : a 6-dim space
defined by 3 spatial coordinates and 3
momentum coordinates

$$\rho_f = \frac{dn}{dE_f} = \frac{dn}{dp} \cdot \frac{dp}{dE} \propto p_e^2$$

Fermi's Golden Rule

The beta-decay must include both the electron and the associated anti-neutrino

$$\frac{d^2n}{dp_e dp_{\bar{\nu}}} \propto p_e^2 p_{\bar{\nu}}^2$$

$$E = E_e + E_{\nu} = E_e + p_{\nu}c$$

$$dp_{\nu} = dE/c$$

$$\therefore \rho \propto p_e^2 (E - E_e)^2$$

Ignores recoil KE of daughter atom.
Assumes neutrino has zero mass

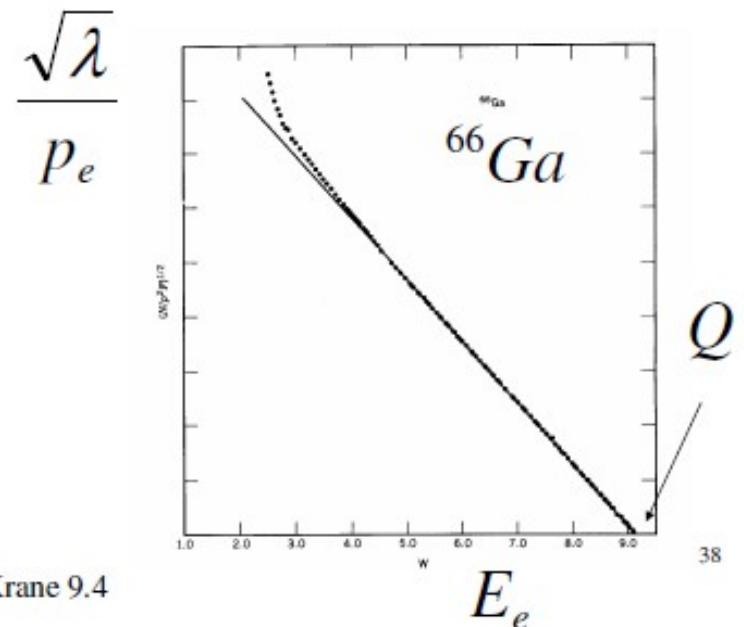
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(Fermi-) Kurie plot

The transition rate is therefore:

$$\lambda \propto p_e^2 (E - E_e)^2$$

$$\therefore \frac{\sqrt{\lambda}}{p_e} \propto (E - E_e)$$



Krane 9.4

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Fermi & Gamow-Teller transitions

Fermi decay: electron and anti-neutrino spins are antiparallel

$$\Delta I = 0 \quad \& \quad \Delta \pi = No$$

Gamow-Teller decay: electron and anti-neutrino spins are parallel

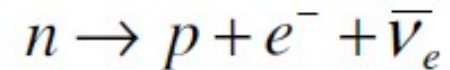
$$\Delta I = 0, \pm 1 \quad \& \quad \Delta \pi = No$$

Except $0 \rightarrow 0$

This assumes that the electron and the anti-neutrino do not carry orbital angular momentum

Calculate the decay rate using the Fermi or the Gamow-Teller mechanisms.

Mirror-decays are mixed transitions:



$$18\% F + 82\% GT$$



Non-Mirror-decays are GT-dominant:

$$\geq 98\% GT$$

β decay is induced by the weak interaction

β -decay mechanism ?

FERMI model assumed a pure VECTOR interaction

$$\Delta I = 0 \quad \& \quad \Delta \pi = No$$

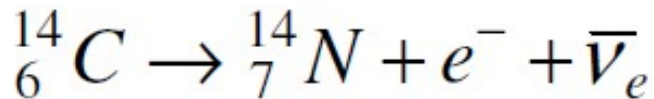
GAMOW-TELLER assumed a TENSOR or AXIAL interaction. Takes into account the spins.

$$\Delta I = 0, \pm 1 \quad \& \quad \Delta \pi = No$$

Now we know that the Nuclear WEAK interaction is Vector-Axial. Strength is about 10^{-6} of the Nuclear Strong force.

*An application of β -decay
"Carbon Dating"*

- The ratio of ^{14}C to ^{12}C is $\sim 1.3 \times 10^{-12}$
- Living organisms exchange CO_2 with surroundings so ratio is \sim stable.
- Exchange stops at death so ratio \downarrow because the ^{14}C decays.



$$t_{1/2} = 5730 \text{ yrs}$$

25 g charcoal

^{14}C activity = 250 decays/min

Decay constant

$$\lambda = \frac{\ln 2}{t_{1/2}} = 3.83 \times 10^{-12} \text{ s}^{-1}$$

Number of ^{12}C nuclei in 25 g

$$= 6.02 \times 10^{23} \times \frac{25}{12} = 1.26 \times 10^{24}$$

$$N_o(^{14}\text{C}) = 1.26 \times 10^{24} \times 1.3 \times 10^{-12} = 1.6 \times 10^{12}$$

Initial activity (at death)

$$R_o = \lambda N_o = 6.13 \text{ decays/sec} = 370 \text{ decays/min}$$

$$R = \lambda N_o e^{-\lambda t}$$

$$250 = 370 e^{-\lambda t} \rightarrow t = 3239 \text{ yrs}$$

D.05. Gamma decay

- Transition between nuclear states:

$$I_i \xrightarrow{\gamma} I_f$$

- A multipole of order L transfers $L\hbar$ angular momentum per photon

$$\vec{I}_i = \vec{L} + \vec{I}_f$$

e.g. $(I_i, I_f) = \left(\frac{3}{2}, \frac{5}{2}\right) \rightarrow L = 1, 2, 3, 4$

i.e. $|I_i - I_f| \leq L \leq (I_i + I_f)$

- ‘Electric’ or ‘Magnetic’ depends on parities of nuclear states

Parity rules for gamma transitions

- Multipole Radiation: Electric and Magnetic
- Opposite parities

$$\pi(EL) = (-1)^L \quad \& \quad \pi(ML) = (-1)^{L+1}$$

- $L = 1 \rightarrow$ Dipole
- $L = 2 \rightarrow$ Quadrupole
- $L = 3 \rightarrow$ Octupole
- $L = 4 \rightarrow$ Hexadecapole etc

- Transition between nuclear states:

$$I_i \xrightarrow{\gamma} I_f$$

- A multipole of order L transfers $L\hbar$ angular momentum per photon

$$\vec{I}_i = \vec{L} + \vec{I}_f$$

e.g. $(I_i, I_f) = \left(\frac{3}{2}, \frac{5}{2}\right) \rightarrow L = 1, 2, 3, 4$

i.e. $|I_i - I_f| \leq L \leq (I_i + I_f)$

- ‘Electric’ or ‘Magnetic’ depends on parities of nuclear states

**Decay operators
in second quantisation:**

gamma transitions **beta transitions**

$$\hat{Y} = \sum_{\tau=p,n} e_{\tau} \sum_{f,i} \langle \tau, f | \hat{V}_{\lambda\mu} | \tau, i \rangle \hat{a}_{\tau f}^+ \hat{a}_{\tau i}$$

$$\hat{\beta}^- = \sum_{f,i} \langle p, f | \hat{V}_{\lambda\mu} | n, i \rangle \hat{a}_{pf}^+ \hat{a}_{ni}$$

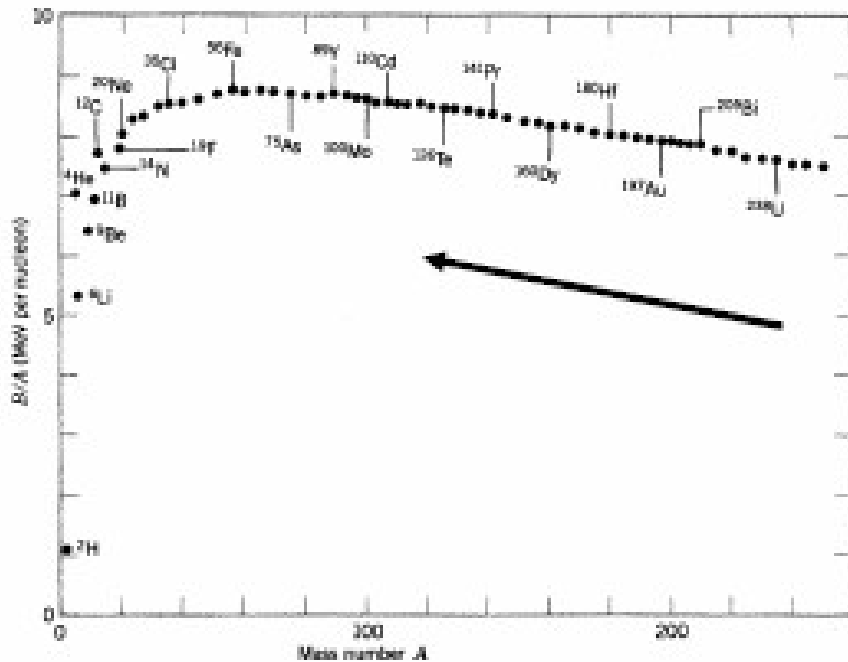
$$\hat{\beta}^+ = \sum_{f,i} \langle n, f | \hat{V}_{\lambda\mu} | p, i \rangle \hat{a}_{nf}^+ \hat{a}_{pi}$$

$$\hat{V}_{\lambda\mu} = r^{\lambda} Y_{\lambda\mu}$$

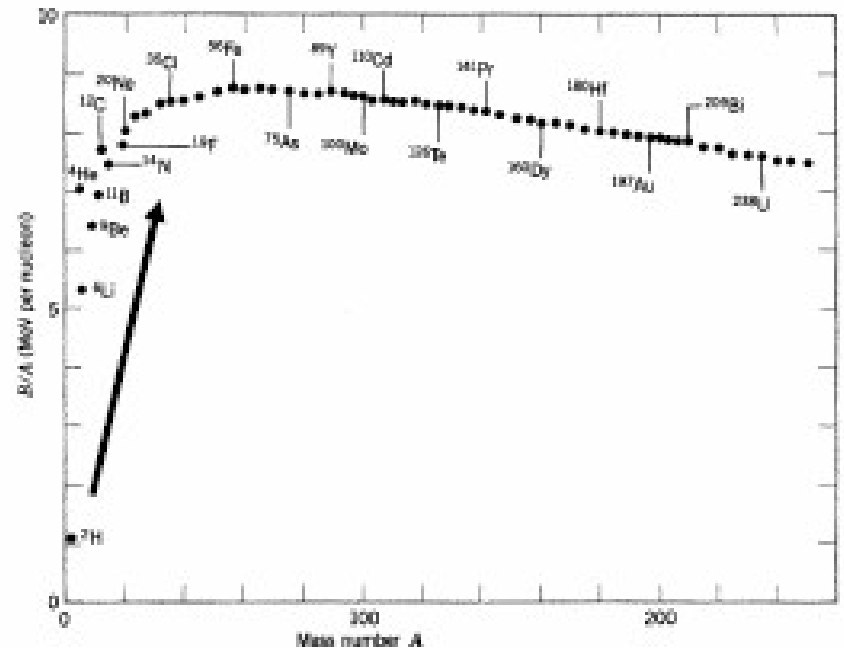
$$\hat{V}_F = 1$$

$$\hat{V}_{GT} = \sigma$$

D.06. Fission & fusion are mirror processes



Fissioning nuclei produce lighter fragments + energy

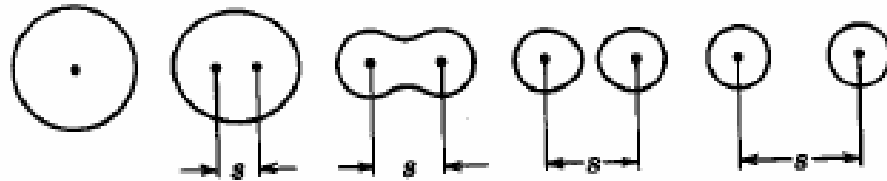


Fusioning nuclei produce heavier fragments + energy

Spontaneous fission liquid drop description

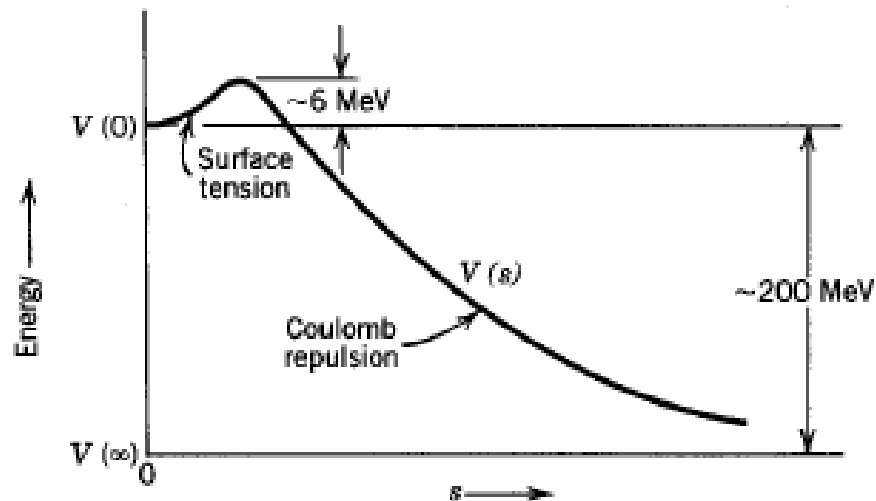
Bohr and Wheeler

Dynamics of the
liquid drop during
the fission process



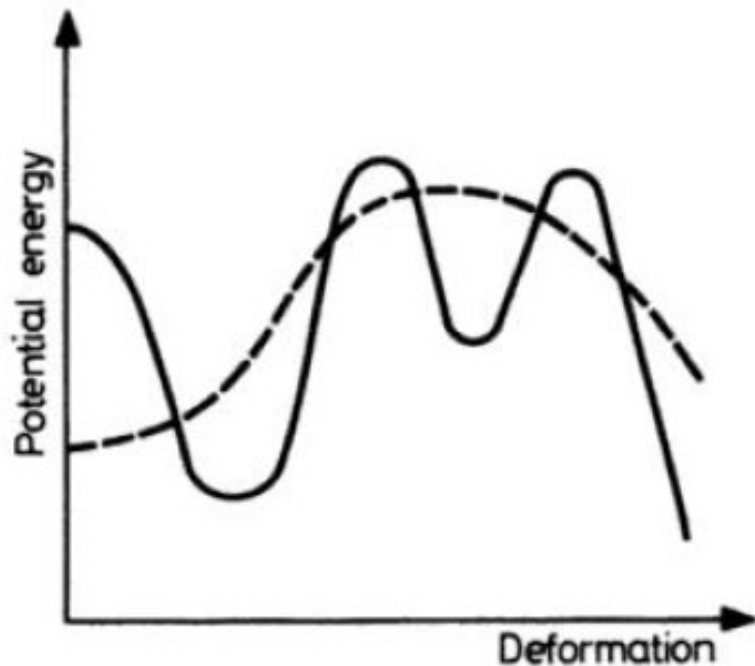
Fission barrier \rightarrow

The decay width is
given by the penetrability
through the Coulomb
barrier (as in the
 α decay theory)



Strutinsky method to compute the fission barrier:

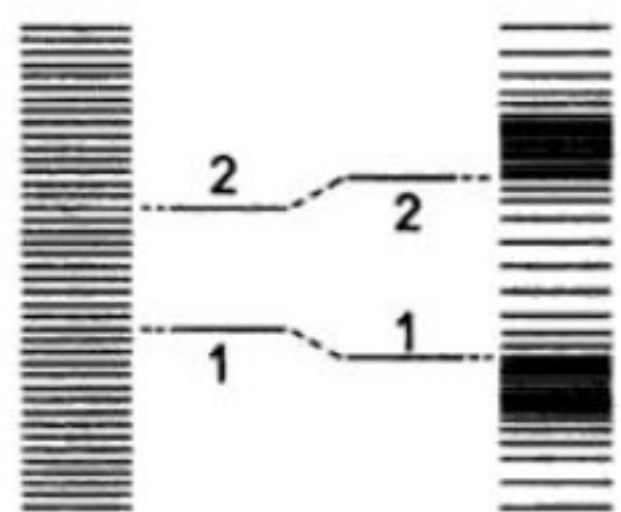
the sum between the liquid drop energy and SM correction
(the difference between SM energy and its mean value).



Density of levels

liquid drop

shell model



Fusion energy

- To fuse two protons, one must overcome the Coulomb repulsion (barrier)
- e.g. two protons 5 fm apart

$$\begin{aligned}U &= \frac{e^2}{4\pi\epsilon_0 r} \\&= (9 \times 10^9)(1.6 \times 10^{-19})^2 / (5 \times 10^{-15}) \\&= 4.6 \times 10^{-14} \text{ J} \quad (288 \text{ keV})\end{aligned}$$

- We have to provide thermal energy of 144 keV to each proton

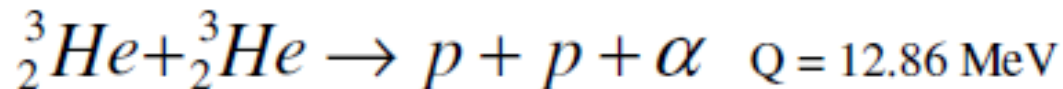
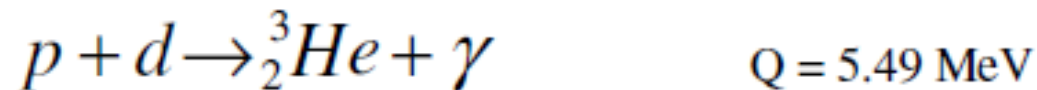
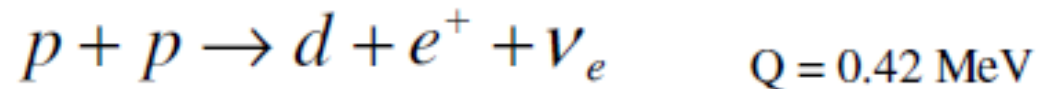
$$\begin{aligned}\langle K \rangle &= \frac{3}{2} k_B T = 2.4 \times 10^{-14} \text{ J} \\ \therefore T &= \frac{2}{3} (2.4 \times 10^{-14} \text{ J}) / (1.381 \times 10^{-23} \text{ J / K}) \\ &= 1.2 \times 10^9 \text{ K}\end{aligned}$$

Energy release for various processes

- Coal (1 kg) → 3×10^7 J
- Oil (1 barrel/0.16 m³) → 6×10^9 J
- Natural Gas (1 m³) → 4×10^7 J
- Uranium (1 kg) → 10^{14} J
- Deuterium fusion (1 kg) → 2×10^{14} J

The Sun

- The temperature of the solar core is about 15 million K.
- The nuclear fusion mechanism in The Sun is the proton-proton cycle.



- The net result is the ‘burning’ of 4 protons to form an alpha particle + positrons, neutrinos, gammas (Eddington).